

RECURSION OPERATOR AND RATIONAL LAX REPRESENTATION

K. Zheltukhin

Department of Mathematics, Faculty of Sciences

Bilkent University, 06533 Ankara, Turkey

Phone: 90 312 2901938 (office)

Fax : 90 312 2664579

e-mail: zhelt@fen.bilkent.edu.tr

Abstract

We consider equations arising from rational Lax representations. A general method to construct recursion operators for such equations is given. Several examples are given, including a degenerate bi-Hamiltonian system with a recursion operator.

PACS: 0230.Ik; 0230.Sr

Keywords: integrable system, recursion operator.

I. Introduction

Recently a new method of constructing a recursion operator from Lax representation was introduced in [1]. This construction depends on Lax representation of a given system of PDEs. Let

$$L_t = [A, L] \quad (1)$$

be Lax representation of an integrable nonlinear system of PDEs. Then a hierarchy of symmetries can be given by

$$L_{t_n} = [A_n, L], \quad n = 0, 1, 2, \dots, \quad (2)$$

where $t_0 = t$, $A_0 = A$ and A_n , $n = 0, 1, 2, \dots$, are Gel'fand-Dikkii operators given in terms of L . The recursion relation between symmetries can be written as

$$L_{t_{n+1}} = LL_{t_n} + [R_n, L], \quad n = 0, 1, 2, \dots, \quad (3)$$

where R_n is an operator such that $\text{ord} R_n = \text{ord} L$.

This symmetry relation allows to find R_n , hence $L_{t_{n+1}}$, in terms of L and L_{t_n} .

In [1], [2] this method was applied to construct recursion operators for Lax equations with different classes of scalar and shift operators, corresponding to field and lattice systems respectively. In [3] the method was applied to Lax equations on a Poisson algebra of Laurent series

$$\Lambda = \left\{ \sum_{-\infty}^{+\infty} u_i p^i : u_i - \text{smooth functions} \right\} \quad (4)$$

with the polynomial Lax function. Such equations give systems of hydrodynamic type. They were also discussed in [4]– [7]. The Hamiltonian structure of the Lax equation on a Poisson algebra was studied in [8].

Here we consider the Lax equation on the Poisson algebra Λ with a rational Lax function

$$L = \frac{\Delta_1}{\Delta_2}, \quad (5)$$

where Δ_1, Δ_2 are polynomials of degree N and M , respectively, and $N > M$. The Lax equation is

$$\frac{\partial L}{\partial t_n} = \{(L)^{\frac{1}{N-M}+n}, L\}, \quad (6)$$

where the Poisson bracket is given by

$$\{f, g\} = p \left(\frac{\partial f}{\partial p} \frac{\partial g}{\partial x} - \frac{\partial f}{\partial x} \frac{\partial g}{\partial p} \right).$$

First we study the symmetry relation (3) for the rational Lax function. Then we give some examples.

In particular, we find a recursion operator \mathcal{R} for equation (6) with the Lax function

$$L = p + S + \frac{P}{p + Q}, \quad (7)$$

which leads to the system [4]

$$\begin{aligned} S_t &= P_x, \\ P_t &= PS_x - QP_x - PQ_x, \\ Q_t &= QS_x - QQ_x. \end{aligned} \quad (8)$$

The recursion operator is given by

$$\mathcal{R} = \begin{pmatrix} S & 1 & PQ^{-1} + P_x D_x^{-1} \cdot Q \\ 2P & S - Q & -2P + (PS_x - (PQ)_x) D_x^{-1} \cdot Q \\ Q & 1 & PQ^{-1} + S - Q + (QS_x - QQ_x) D_x^{-1} \cdot Q \end{pmatrix} \quad (9)$$

In [4] bi-Hamiltonian representation of this equation was constructed with Hamiltonian operators

$$\mathcal{D}_1 = \begin{pmatrix} 0 & P & Q \\ P & -2PQ & -Q^2 \\ Q & -Q^2 & 0 \end{pmatrix} D_x + \begin{pmatrix} 0 & P_x & Q_x \\ 0 & -(PQ)_x & -QQ_x \\ 0 & -QQ_x & 0 \end{pmatrix} \quad (10)$$

and

$$\mathcal{D}_2 = \begin{pmatrix} 2P & P(S - 3Q) & Q(S - Q) \\ P(S - 3Q) & P(2P - 2SQ + 4Q^2) & Q(2P - SQ + Q^2) \\ Q(S - Q) & Q(2P - SQ + Q^2) & 2Q^2 \end{pmatrix} D_x + \quad (11)$$

$$\begin{pmatrix} P_x & SP_x - 2(PQ)_x & SQ_x - QQ_x \\ PS_x - (QP)_x & (-SPQ + P^2 + 2PQ^2)_x & Q_x(2P + Q^2 - SQ) \\ QS_x - QQ_x & Q(2P_x + 2QQ_x - S_x - SQQ_x) & 2QQ_x \end{pmatrix}$$

These Hamiltonian operators are degenerate, so, one can not use them to find a recursion operator. But it turns out that they are related to the recursion operator \mathcal{R} . One can easily check that the following equality holds

$$\mathcal{R}\mathcal{D}_1 = \mathcal{D}_2.$$

We observe that the degeneracy in the bi-Hamiltonian operators is due to the following fact. Let $p' = p + F$ then the Lax function becomes

$$L = p' + G + \frac{P}{p'} . \quad (12)$$

This means that we have two independent variabls P and G , where $G = S - F$. The equation corresponding to the Lax function (12) has been studied in [3].

To remove degeneracy one can take the Lax function as

$$L = p + S + \frac{P}{p} + \sum_{i=1}^m \frac{Q_i}{p + F_i} . \quad (13)$$

As an example we shall consider the equation (6) with the Lax function

$$L = p + S + \frac{P}{p} + \frac{Q}{p + F} . \quad (14)$$

II. Symmetry Relation for Rational Lax Representation.

Following [1] we consider the hierarchy of symmetries for the Lax equation (6) with the Lax function (5)

$$\frac{\partial L}{\partial t_n} = \{(L^{\frac{1}{N-M}+n})_{\geq 0}, L\} . \quad (15)$$

Lemma 1. *For any $n = 0, 1, 2, \dots$,*

$$\frac{\partial L}{\partial t_n} = L \frac{\partial L}{\partial t_{n-1}} + \{R_n, L\} . \quad (16)$$

Function R_n has a form

$$R_n = A + \frac{B}{\Delta_2} \quad (17)$$

where A is a polynomial of degree $(N - M)$ and B is a polynomial of degree $(M - 1)$.

Proof. We have

$$(L^{\frac{1}{N-M}+n})_{\geq 0} = [L(L^{\frac{1}{N-M}+(n-1)})_{\geq 0} + L(L^{\frac{1}{N-M}+(n-1)})_{< 0}]_{\geq 0}$$

So,

$$\begin{aligned} (L^{\frac{1}{N-M}+n})_{\geq 0} &= L(L^{\frac{1}{N-M}+(n-1)})_{\geq 0} + (L(L^{\frac{1}{N-M}+(n-1)})_{< 0})_{\geq 0} - \\ &\quad (L(L^{\frac{1}{N-M}+(n-1)})_{\geq 0})_{< 0}. \end{aligned}$$

If we take

$$R_n = (L(L^{\frac{1}{N-M}+(n-1)})_{< 0})_{\geq 0} - (L(L^{\frac{1}{N-M}+(n-1)})_{\geq 0})_{< 0} \quad , \quad (18)$$

then

$$(L^{\frac{1}{N-M}+n})_{\geq 0} = L(L^{\frac{1}{N-M}+(n-1)})_{\geq 0} + R_n.$$

Hence,

$$\frac{\partial L}{\partial t_n} = \left\{ (L^{\frac{1}{N-M}+n})_{\geq 0}; L \right\} = \left\{ L(L^{\frac{1}{N-M}+(n-1)})_{\geq 0} + R_n; L \right\} = L \frac{\partial L}{\partial t_n} + \{R_n; L\},$$

and (16) is satisfied. The remainder R_n has form (17). Indeed in (18) we set

$$A = (L(L^{\frac{1}{N-M}+(n-1)})_{< 0})_{\geq 0}$$

and

$$B = \Delta_2 \cdot (L(L^{\frac{1}{N-M}+(n-1)})_{\geq 0})_{< 0}$$

Then A is a polynomial of degree $(N - M - 1)$ and B is a polynomial of degree $(M - 1)$. \square

Now we can apply the Lemma 1 to find recursion operators.

III. Examples.

Example 2. Let us consider the equation (8) given in introduction.

Lemma 3. *A recursion operator for (8) is given by (9).*

Proof. Using (17) for R_n , we have $R_n = A + \frac{B}{p + Q}$. So, the symmetry relation (16) is

$$\begin{aligned} & \frac{\partial S}{\partial t_n} + \frac{\partial P}{\partial t_n} \cdot \frac{1}{p + Q} + \frac{\partial Q}{\partial t_n} \cdot \frac{P}{(p + Q)^2} = \\ & \left(p + S + \frac{P}{p + Q} \right) \left(\frac{\partial S}{\partial t_{n-1}} + \frac{\partial P}{\partial t_{n-1}} \cdot \frac{1}{p + Q} + \frac{\partial Q}{\partial t_{n-1}} \cdot \frac{P}{(p + Q)^2} \right) + \\ & p \left(A_x + \frac{B_x}{p + Q} + \frac{-BQ_x}{(p + Q)^2} \right) \left(1 + \frac{-P}{(p + Q)^2} \right) \\ & - \frac{pB}{(p + Q)^2} \left(S_x + \frac{P_x}{p + Q} + \frac{-PQ_x}{(p + Q)^2} \right) \end{aligned}$$

To have the equality the coefficients of p and $(p + Q)^{-3}$ must be zero. It gives the recursion relations to find A and B . Then the coefficients of p^0 , $(p + Q)^{-1}$, $(p + Q)^{-2}$ give expressions for $\frac{\partial S}{\partial t_n}$, $\frac{\partial P}{\partial t_n}$, $\frac{\partial Q}{\partial t_n}$. \square

Example 4. The Lax equation (6) with the Lax function (14), for $n = 1$, gives the following system

$$\begin{aligned}
S_t &= P_x + Q_x, \\
P_t &= PS_x, \\
Q_t &= QS_x - FQ_x - QF_x, \\
F_t &= FS_x - FF_x.
\end{aligned} \tag{19}$$

Lemma 5. *A recursion operator for (19) is given by*

$$\begin{pmatrix}
S & 2 + P_x D_x^{-1} \cdot P^{-1} & 1 & QF^{-1} + Q_x D_x^{-1} \cdot F^{-1} \\
2P & S + QF^{-1} + PS_x D_x^{-1} \cdot P^{-1} & PF^{-1} & -2PQF^{-2} \\
& PF^{-1}(Q_x - QF^{-1}F_x)D_x^{-1} \cdot P^{-1} & & -PF^{-1}(Q_x - QF^{-1}F_x)D_x^{-1} \cdot F^{-1} \\
2Q & -QF^{-1} & S - F & -2PQF^{-2} - 2Q \\
& -PF^{-1}(Q_x - QF^{-1}F_x)D_x^{-1} \cdot P^{-1} & -PF^{-1} & +PF^{-1}(Q_x - QF^{-1}F_x)D_x^{-1} \cdot F^{-1} \\
& & & +(QS_x - QF_x - FQ_x)D_x^{-1} \cdot F^{-1} \\
F & 1 + (P_x - PF^{-1}F_x)D_x^{-1} \cdot P^{-1} & -1 & PF^{-1} - F + (FS_x - FF_x)D_x^{-1} \cdot F^{-1} \\
& & & -(P_x - PF^{-1}F_x)D_x^{-1} \cdot F^{-1}
\end{pmatrix}. \tag{20}$$

Proof. Using (17) for R_n , we have $R_n = C + \frac{A}{p} + \frac{B}{p+F}$. So, the symmetry relation (16) is

$$\begin{aligned}
& \frac{\partial S}{\partial t_n} + \frac{\partial P}{\partial t_n} \cdot \frac{1}{p} + \frac{\partial Q}{\partial t_n} \cdot \frac{1}{(p+F)} + \frac{\partial F}{\partial t_n} \cdot \frac{-Q}{(p+F)^2} = \\
& \left(p + S + \frac{P}{p} + \frac{Q}{p+F} \right) \left(\frac{\partial S}{\partial t_{n-1}} + \frac{\partial P}{\partial t_{n-1}} \cdot \frac{1}{p} + \frac{\partial Q}{\partial t_{n-1}} \cdot \frac{1}{(p+F)} + \frac{\partial F}{\partial t_{n-1}} \cdot \frac{-Q}{(p+F)^2} \right) +
\end{aligned}$$

$$p \left(\frac{-B}{p^2} + \frac{-C}{(p+F)^2} \right) \left(S_x + \frac{P_x}{p} + \frac{Q_x}{(p+F)} + \frac{-QF_x}{(p+F)^2} \right) -$$

$$p \left(A_x + \frac{B_x}{p} + \frac{C_x}{(p+F) + \frac{-CF_x}{(p+F)^2}} \right) \left(1 + \frac{P}{p} + \frac{-Q}{(p+F)^2} \right)$$

Therefore, the coefficients of p , p^{-2} , and $(p+F)^{-3}$ must be zero, it gives recursion relations to find A , B and C . Then the coefficients of p^0 , p^{-1} , $(p+F)^{-1}$ and $(p+F)^{-2}$, give expressions for $\frac{\partial S}{\partial t_n}$, $\frac{\partial P}{\partial t_n}$, $\frac{\partial Q}{\partial t_n}$ and $\frac{\partial F}{\partial t_n}$. \square

Acknowledgments

I thank Professors Metin Gürses, Atalay Karasu and Maxim Pavlov for several discussions. This work is partially supported by the Scientific and Technical Research Council of Turkey.

References

- [1] M. Gürses, A. Karasu, V.V. Sokolov, *J. Math. Phys.*, **40**, 6473-6490 (1999).
- [2] M.Blaszak ” On construction of recursion operator and algebra of symmetries for field and lattice systems”, to appear in *Rep. Math. Phys.*
- [3] M. Gürses and K. Zhelthukin, *J. Math. Phys.* **42**, 1309-1325 (2001).
- [4] I.A.B. Strachan, *J. Math. Phys.*, **40**, 5058-5079 (1999).
- [5] D.B. Fairlie and I.A.B. Strachan, *Inverse Problems*, **12**, 885-908 (1998).

- [6] J.C. Brunelli, M. Gürses, and K. Zhelthukin, *Reviews in Mathematical Physics*, vol. 13, No. 4, 529-543 (2001).
- [7] J.C. Brunelli and A. Das, *Phys. Lett. A* ,235 , 597-602 (1997).
- [8] Luen-Chau Li, *Commun. Math Phys.*, 203, 573-592 (1999).